

a) Define a Fourier transform in the t -direction,

$$\mathcal{F}_t[u(x, t)] = \int_{-\infty}^{\infty} u(x, t) e^{-i\omega t} dt = U(x, \omega)$$

Transforming the PDE in the t -direction,

$$\begin{aligned} \mathcal{F}_t\left[\frac{\partial^2 u}{\partial x^2} + 2\frac{\partial u}{\partial t} + \alpha u\right] &= \mathcal{F}[0] \\ \Rightarrow \mathcal{F}_t\left[\frac{\partial^2 u}{\partial x^2}\right] + 2\mathcal{F}_t\left[\frac{\partial u}{\partial t}\right] + \alpha \mathcal{F}[u] &= 0 \text{ by linearity} \\ \Rightarrow \frac{\partial^2}{\partial x^2} U(x, \omega) + 2i\omega U(x, \omega) + \alpha U(x, \omega) &= 0 \text{ by Theorems 16.1 and 16.2} \\ \Rightarrow \frac{\partial^2}{\partial x^2} U(x, \omega) + (2i\omega + \alpha) U(x, \omega) &= 0 \end{aligned}$$

The auxiliary equation is

$$m^2 + 2i\omega + \alpha = 0 \Rightarrow m = -2i\omega - \alpha \text{ so}$$

$$U(x, \omega) = A(\omega) e^{-(2i\omega + \alpha)x} \quad (1)$$

Transforming the boundary condition to find $A(\omega)$,

$$\mathcal{F}_t[u(0, t)] = \mathcal{F}_t[f(t)] \Rightarrow U(0, \omega) = F(\omega)$$

where $F(\omega) = \mathcal{F}[f(t)]$.

From (1),

$$U(0, t) = A(\omega) e^0 = A(\omega)$$

Hence

$$A(\omega) = F(\omega) \text{ so}$$

$$U(x, \omega) = F(\omega) e^{-(2i\omega + \alpha)x}$$

Then define $G(x, \omega) = e^{-(2i\omega + \alpha)x}$ so

$$U(x, \omega) = F(\omega) G(x, \omega)$$

b) From the standard transforms,

$$\mathcal{F}^{-1}[e^{-i\omega\tau}] = \delta(t - \tau)$$

$$G(x, \omega) = e^{-2i\omega x - \alpha x} = e^{-\alpha x} e^{-2i\omega x}$$

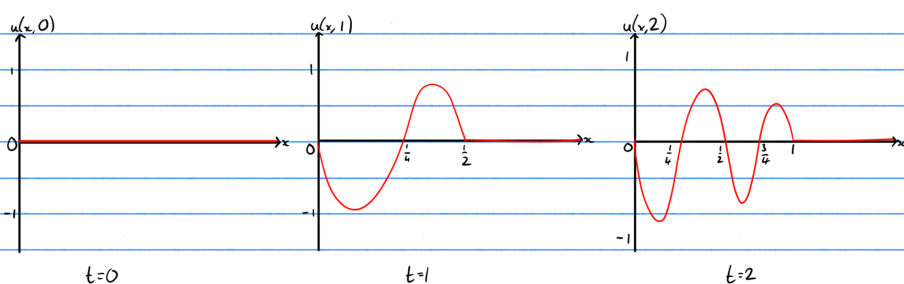
$e^{-\alpha x}$ only depends on x hence with $\tau = 2x$ in the standard transform,

$$g(x, t) = \mathcal{F}_t^{-1}[G(x, \omega)] = e^{-\alpha x} \mathcal{F}_t^{-1}[e^{-2i\omega x}] = e^{-\alpha x} \delta(t - 2x) \text{ by linearity}$$

c) $\mathcal{F}_t^{-1}[U(x, \omega)] = \mathcal{F}_t^{-1}[F(\omega) G(x, \omega)]$

$$\begin{aligned} \Rightarrow u(x, t) &= (f * g)(x, t) \\ &= \int_{-\infty}^{\infty} f(\tau) g(x, t - \tau) d\tau \\ &= e^{-\alpha x} \int_{-\infty}^{\infty} f(\tau) \delta(t - 2x - \tau) d\tau \\ &= e^{-\alpha x} f(t - 2x) \text{ by the sifting property of the delta function.} \end{aligned}$$

$$d) u(x, t) = \begin{cases} e^{-\alpha x} H(t - 2x) \sin(2\pi(t - 2x)) & x \geq 0 \\ e^{-\alpha x} \sin(2\pi(t - 2x)) & x < \frac{1}{2}t \\ 0 & x \geq \frac{1}{2}t \end{cases}$$



e) Up to $x = \frac{1}{2}t$, the solutions exhibit sinusoidal behaviour which reduces in amplitude as x increases. Post $x = \frac{1}{2}t$ the solution is 0. A physical system might be sound waves being produced and propagating through a medium where $x=0$ is the point where the sound is produced.